

Introduction to the study of surface by Small Angle X-ray Scattering









Thomas ZEMB Diane REBISCOUL





IIG





Absolute intensity





Porod's law



If q >> curvature \Rightarrow flat at q⁻¹ scale \Rightarrow **Porod's Regime** $d = 2\pi/q$ **Porod's law** $I(q) = 2\pi(\rho_1 - \rho_2)^2 \Sigma \frac{1}{q^4}$ Total interfacial area per unit volume (S/V)



Thus Porod's law traduces the reflection of the incident beam on the portions of interface in the samples which have the right orientation to yield a reflected beam at a given scattering angle. In this sense, Porod's law is nothing else than Fresnel's law of reflectivity calculated in the Born approximation, a point of view particularly developed in ref. [7].

L. Auvray and P. Aurroy, Bombannes summer school lecture : Scatering by interfaces (1991)





Effect of curvature



 $q^{-1} \equiv$ Interface curvature

Kirste-Porod Formula => addition of q⁻⁶ term

$$I(q) = 2\pi(\rho_1 - \rho_2)^2 \sum_{q=1}^{1} (1 + \frac{1}{q^2} \left[\frac{1}{4} < (C_1 + C_2)^2 > + \frac{1}{8} < (C_1 - C_2)^2 > \right]$$

$$Local principal surface curvature$$

$$Kirste - Porod$$

$$Correction$$

$$Porod's plateau$$

L. Auvray and P. Aurroy, Bombannes summer school lecture : Scaterring by interfaces (1991) R. Kirste, G. Porod, Koll, Z., Z. für Polym, 184, 1 (1962)



When the interface is grafted or covered



L. Auvray et al , Microemulsions: a small angle X-ray scattering study. Journal de Physique, 1984, 45 (5), pp.913-928. L. Auvray, J.P. Cotton, R. Ober, C. Taupin, Structure of concentrated Windsor microemulsion by SANS, Physica 136B (1986) 281-283.



Diffuse interface

1/E Smoothing (Ghost) Function, h(r) 0.606 negative deviations from the Porod law due to the diffuse interface Expected Electron Density Profile, $\tilde{n}(r) = g(r) * h(r)$ Modified Porod's Law $I_{FICor}q^{4} = K_{p} \frac{\sin^{2}(E \cdot q/2)}{(E \cdot q/2)^{2}} \qquad I_{FICor}q^{4} = K_{p} \left[\frac{4\sin(T \cdot q/4)}{(T \cdot q)}\right] I_{FICor}q^{4} = K_{p} \cdot \exp(-\sigma^{2} \cdot q^{2})$ $\lim_{q\to\infty} [I(q)] = I_p(q)H^2(q) + I_{FI}(q)$ - exact $I_{\text{PICor}}q^4 = K_p(1 - \frac{E^2 \cdot q^2}{12})$ $I_{\text{PICor}}q^4 = K_p(1 - \frac{T^2 \cdot q^2}{24})$ expanded $I_{\text{PlCor}}q^4 = K_{n} \cdot (1 \cdot \sigma^2 \cdot q^2)$ **Density fluctuation** (for $\sigma \cdot q \ll 1.0$) (for $E \cdot q/2 \ll 1.0$) (for T · q/4 << 1.0) within the diffuse layer Required Plot exact $\ln(I_{FCor}q^4)$ vs. q^2 I FICor q4 vs. q2 I FICor q4 vs. q2 IFICerq⁴ vs.q² expanded $I_{FIcor} = I_p(q)H^2(q)$ Determination of transition layer thickness, E and T standard deviation, o $K_p = 2\pi \Sigma (\Delta \rho)^2$ - exact $\sigma = \sqrt{-m}$ - expanded^{a)} $E = \sqrt{\frac{-12 \cdot m}{K_p}}$ K_p = Porod constant $\sigma = \sqrt{-m/K_o}$ $T = 2\sqrt{\frac{-6 \cdot m}{V}}$ $T(=2\sqrt{6\sigma}) > E(=\sqrt{12\sigma}) > \sigma$ Relationship Vonk (1973) Ruland (1971) References this work Kim, J. of Appl. Cryst, 2003, 643-651 (a) K_p and m are Porod's constant (see the text for definition) and slope in each plot.



Diffuse interface

Modified Porod's law plot for sigmoidal (σ model) gradient

$$I_{FiCor}q^4 = K_p \cdot exp(-\sigma^2 \cdot q^2)$$



$$I_{PICor}q^{4} = K_{p} \cdot (1 - \sigma^{2} \cdot q^{2})$$

$$I_{Obs}(q) \qquad (for \sigma \cdot q << 1.0)$$

$$I_{FI}(q)$$

$$K_{p}=0.1271$$

$$K_p = 2\pi \Sigma (\Delta \rho)^2$$

Crystallized ethylene/1-octene (EO) random copolymers with hexyl branch content

Heterogenous media at all lengthscale

Fully random dispersion of 2 phases at any scale => Debye Model



$$I(q) \sim \frac{1}{(1+q^2 l_p^2)^2} \longrightarrow l_p \Sigma = 4\Phi(1-\Phi)$$
Porod's lenght
Volume fraction
of one phase

 \odot $I(q)q^4 =>$ monotonously increasing + no asymptotic behavior

⊗ Not satisfying for smooth surfaces

⊗ Not realistic since it assumes a too broad of pore sizes

P. Debye, H.R. Henderson, H. Brumberger, J. Appl. Phys. 28, 679 (1957) E.W. Käler, S. Prager, J. Coll. Interf. Sci. 86, 359 (1982).





D < 3 Diffusion from volume

More complicated

Wong et al, 1988

Length scale

 $R = 4 \text{ Å for } N_2$



Mass and surface fractal structures

Several instruments to cover a large q range



A. P. Radlinski, E. Z. Radlinsk , M. Agamalian , G. D. Wignall , P. Lindner and O. G. Randl, The fractal microstructure of ancient sedimentary rocks, J. Appl. Cryst. (2000) 33, 860-862.

Seneralization to packing of polydisperse particles

Spherical primary particles



Beaucage 1995, 1996, 2004





Highly absorbing materials

 UO_2 or $ThO_2 => Absorption length \approx 10 \ \mu m (17.4 \ keV)$





Highly absorbing materials

How to obtain the I abs with the right transmission?

1. Measured transmission T^{mes} as a function of the vertical distance z from the interface powder/air





Highly absorbing materials

How to obtain the I abs with the right transmission

2. I_S^{mes} has to be corrected from the multi-scattering and indirect absorption effects occurring in low transmission case



Highly absorbing materials

How to obtain the I abs with the right transmission



$$I_{S}^{abs} = \frac{1}{J_{in}A\epsilon\Delta\Omega e_{b}} \left(\frac{\alpha I_{S}^{mes}(sample) - I_{B}}{t_{sample}T_{sample}^{mes}} - \frac{I_{S}^{mes}(tube) - I_{B}}{t_{tube}T_{tube}^{mes}}\right)$$
$$\frac{S}{V} = \frac{\lim_{Q \to \infty} I_{S}^{abs} q^{4}}{2\pi\Delta SLD^{2}}$$

ESRF (16 keV) => lower absorption and smaller beam size

position	S (m²/g)	Porosity ϕ
2	60	0.08
3	80	0.11
4	127	0.19
5	284	0.34
ESRF	310	0.45